

PHYSICALLY-BASED CONSTITUTIVE EQUATIONS FOR IRRADIATED RPV AND INTERNALS

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- ❑ What is a crystalline law ?
- ❑ The rate equation
- ❑ Stress components
- ❑ Microstructure evolution during deformation

What is a crystalline law ?



The stress-strain curve is needed in many mechanical applications

Non-crystallographic Laws

$$\sigma = \sigma_o + k \varepsilon^n$$

[Hollomon, Ludwig]

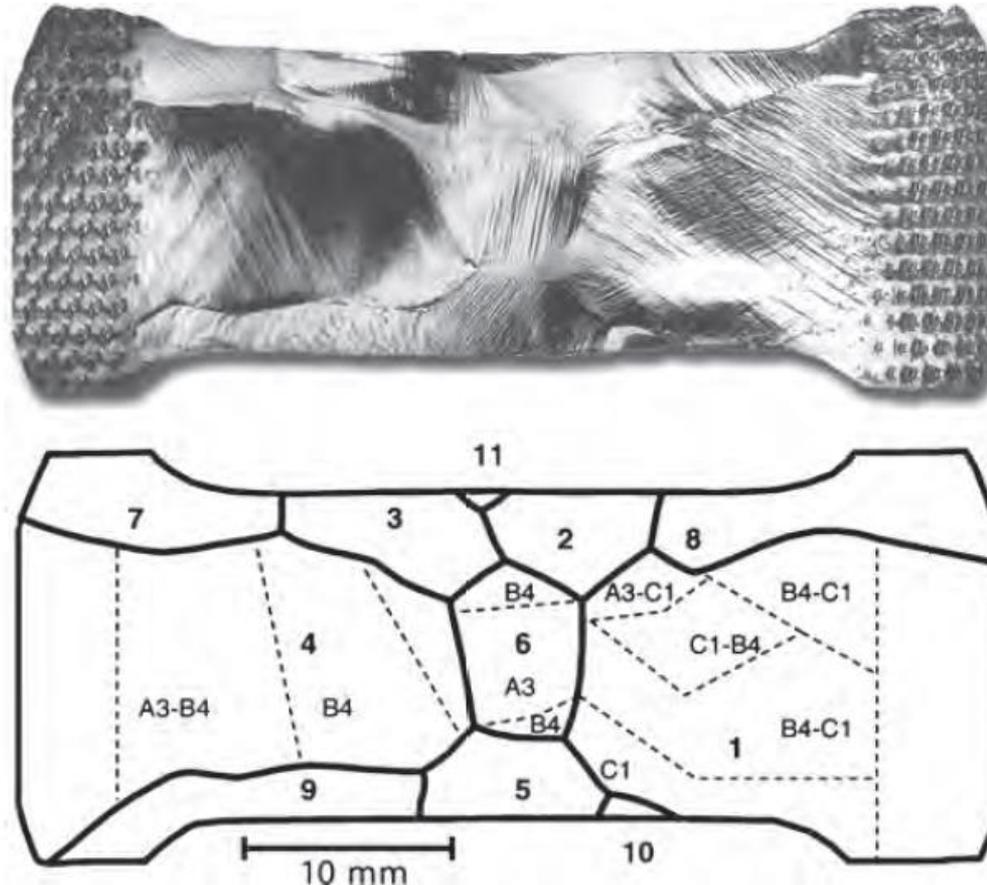
$$\frac{\varepsilon}{\varepsilon_o} = \frac{\sigma}{\sigma_o} + \alpha \left(\frac{\sigma}{\sigma_o} \right)^n$$

$\sigma_o = E\varepsilon_o$ Ramberg-Osgood

- ❑ Not predictive
- ❑ Uniform homogeneous local/macroscopic behavior
- ❑ No explicit grain representation, nor physical properties

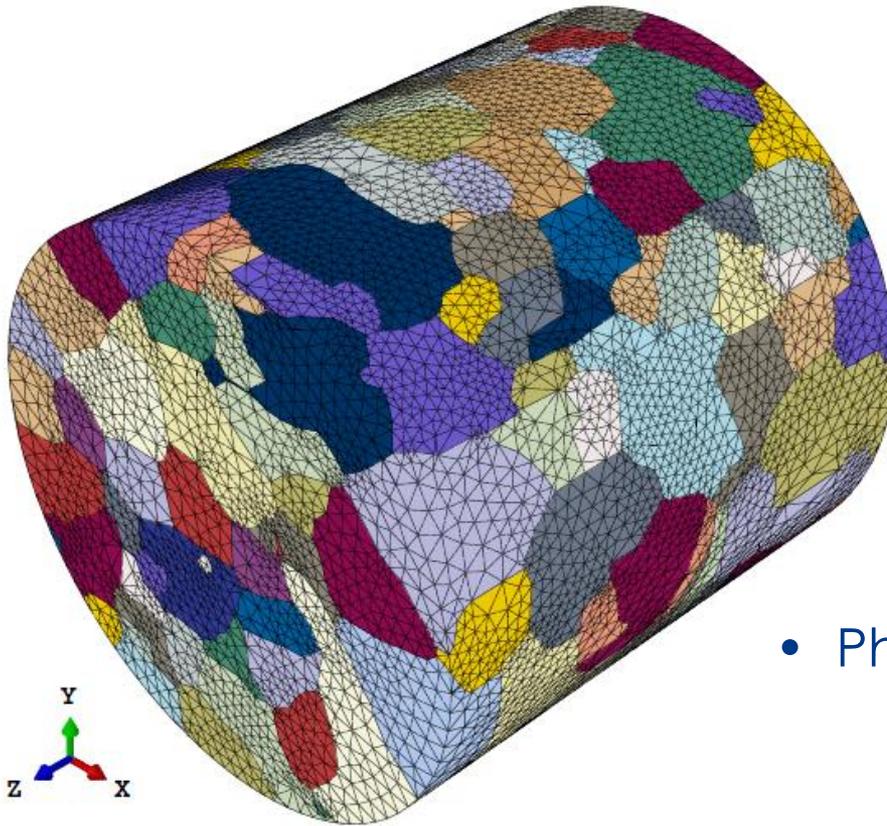
F. Roters et al. Crystal Plasticity Finite Element Methods:
in Materials Science and Engineering, Wiley VCH, Weinheim, 2010.

What is a crystalline law ?



In metallic alloys, plasticity is accommodated by slip on specific planes along specific directions

What is a crystalline law ?



Homogenization calculations

+ Boundary conditions

→ Local stress field

Need mechanical local response

- Phenomenological crystalline Laws

$$\tau_c^s = f(T, \gamma_{cum}^i)$$

No microstructure explicit variables

What is a crystalline law ?



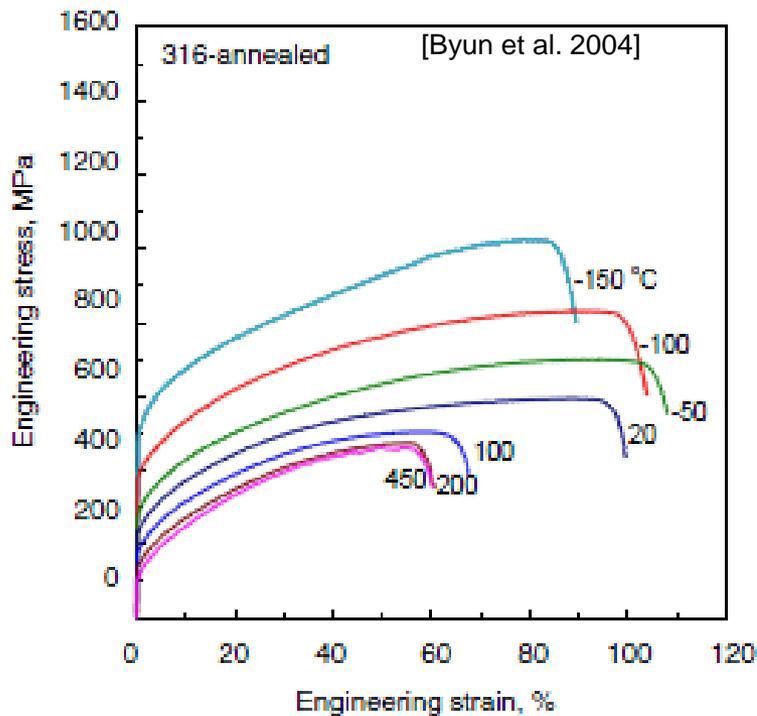
□ Physically-based Crystalline Laws

- Mechanical properties are determined by **physical properties** P_i (elastic constants, grain size, dislocation density, obstacles size/density, etc.) AND **test parameters** T_i (temperature, strain rate, applied load, etc.)

- The rate equation $\dot{\epsilon} = f(P_i, \tau_{eff}^s)$ $\tau_{eff}^s ?$
- Stress decomposition $\tau_{app}^s = f(\tau_f^s, \tau_{HP}^s, \tau_{forest}^s, \tau_{sc}^s, \tau_{eff}^s \dots)$
- Microstructure evolution $\partial Q_i^s = f(T_i, P_i, \gamma^s, \Delta\gamma)$

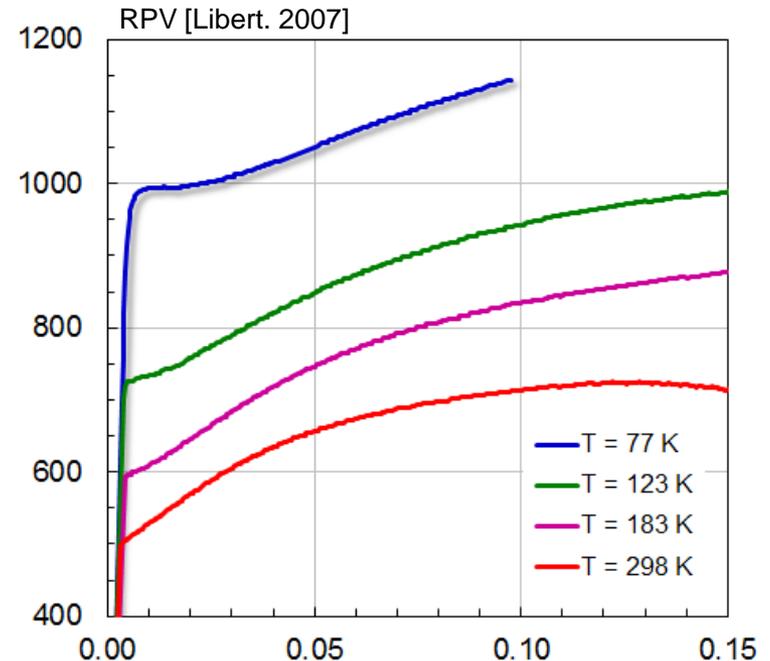
What is a crystalline law ?

FCC – type behavior



- unambiguous slip planes
- no fragile regime
- temperature-dependent parameters

BCC – type behavior

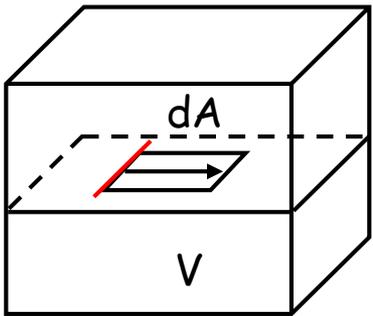


- many slip planes/non-crystallographic
- ductile-fragile transition
- Temperature-dependent mechanisms

- ❑ What is a crystalline law ?
- ❑ The rate equation
- ❑ Stress components
- ❑ Microstructure evolution during deformation

Rate equation: the Orowan relation

If a dislocation of length L and Burgers vector b moves by dx in a volume V :



$$d\gamma = \frac{bdA}{V} = \frac{bLdx}{V} = b\rho dx$$

$$\dot{\gamma} = \frac{d\gamma}{dt} = \frac{b\rho dx}{dt} = b\rho v$$

$$\dot{\gamma} = b\rho_m^s v^s$$

Orowan relation

$$\tau_{app}^s = \tau_c^s + \tau_{eff}^s$$

Dislocation velocity and strain rate = $f(\tau_{eff})$

Rate equation

$$\dot{\gamma} = b\rho_m^s v_o \exp\left(-\frac{\Delta G(\tau_{eff}^s)}{kT}\right)$$

Rate equation: a thermal regime



$$\dot{\gamma}^s = b \rho_m^s v_o \exp\left(-\frac{\Delta G(\tau_{eff}^s)}{kT}\right) \quad V^* = -\frac{\partial \Delta G}{\partial \tau_{eff}^s}, \quad \beta = \frac{1}{kT}$$

Stress components $\tau_c^s = \lim_{\dot{\gamma}^s \rightarrow 0} (\tau_{app}^s) \quad \tau_{app}^s = \tau_{eff}^s + \tau_c^s$

When $\tau_{eff}^s \ll \tau_c^s$ (FCC – type behavior)

$$\dot{\gamma}^s = b \rho_m^s v_o \exp[-\beta(A - V\tau_{eff}^s)] = \dot{\gamma}_o^s \exp(\beta V \tau_{eff}^s) = \dot{\gamma}_o^s \left[\exp\left(\frac{\tau_{app}^s}{\tau_c^s} - 1\right) \right]^{\beta V \tau_c^s}$$

$$\dot{\gamma}^s = \dot{\gamma}_o^s \left(\frac{\tau_{app}^s}{\tau_c^s} \right)^{\beta V \tau_c^s} = \dot{\gamma}_o^s \left(\frac{\tau_{app}^s}{\tau_c^s} \right)^n$$

Rate equation: a thermal regime

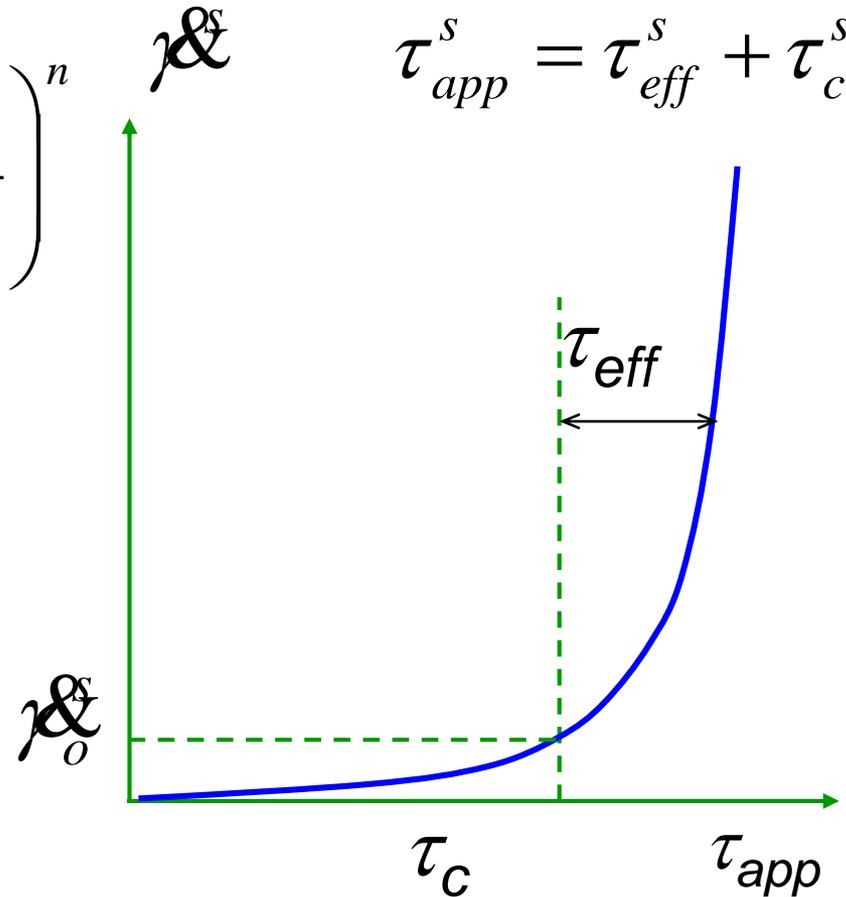


$$\dot{\epsilon} = \dot{\epsilon}_0 \left(\frac{\tau_{app}^s}{\tau_c^s} \right)^{BV\tau_c^s} = \dot{\epsilon}_0 \left(\frac{\tau_{app}^s}{\tau_c^s} \right)^n$$

$$\tau_{app}^s = \tau_{eff}^s + \tau_c^s$$

$k = 1.4E-24$
 $T = 300 \text{ K}$
 $V = 100 \text{ b}^3$
 $\tau_c = 50 \text{ MPa}$

n = 190



τ_c plays the role of threshold of the flow stress

- ❑ What is a crystalline law ?
- ❑ The rate equation
- ❑ Stress components
- ❑ Microstructure evolution during deformation

- The applied stress must exceed:
 - Solid solution: τ_f
 - Hall-Petch (grain size) effect: τ_{HP}
 - Local obstacles (dislocations, precipitates, loops..): τ_{obs}

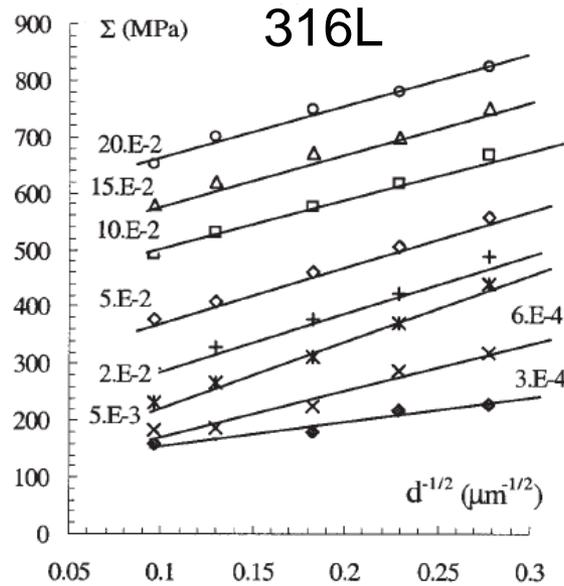
- To overcome lattice and phonon dynamical friction : the effective stress τ_{eff}

Stress components: friction stress τ_f



- ❑ Stress necessary to start moving dislocations
- ❑ Results from:
 - interstitial solute atoms (solid solution of C, N , H, etc.)
 - substitutional solute atoms: Ni in Al, Cu in Al, (Cr, Mn, Ni, S, Mo, etc.) in Fe
- ❑ Depends on temperature and strain rate
- ❑ Cannot be predicted (no reliable model)
- ❑ Material property

Stress components: Hall-Petch effect τ_{HP}



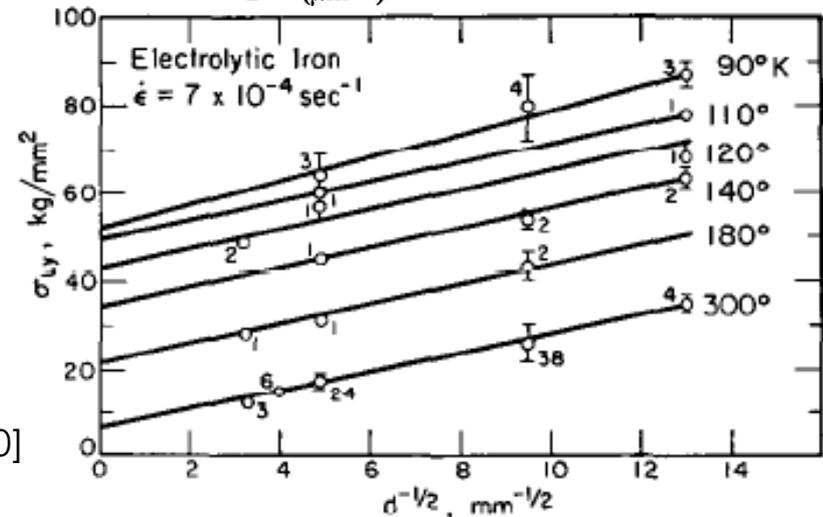
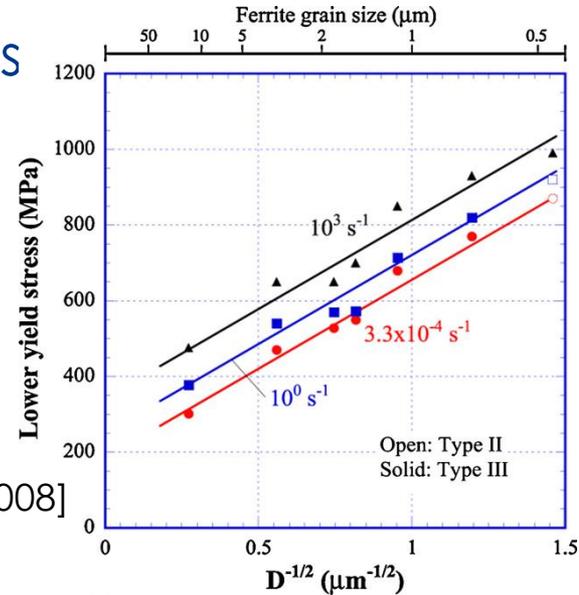
[Feaugas et al. 2003]

$$\tau_{HP} = \frac{\mu}{\mu(300 K)} \frac{K}{\sqrt{d_{grain}}}$$

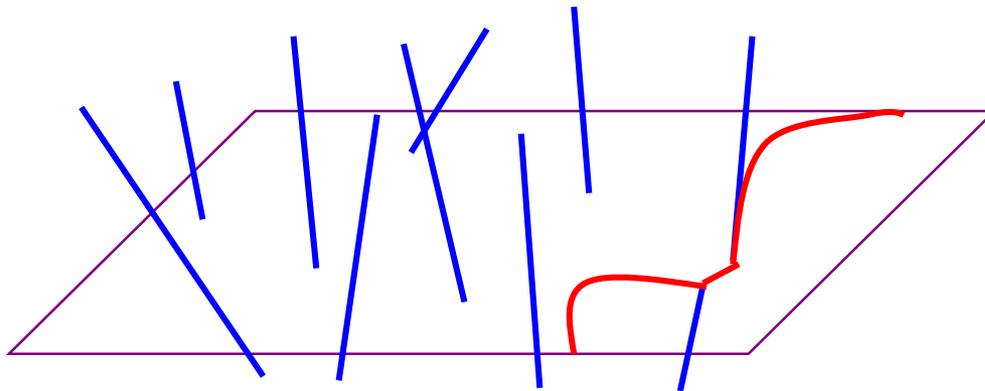
[Conrad et al. 1960]

Ferritic steels

[Tsushida et al 2008]



- ❑ Principal obstacles in all materials
- ❑ Resulting from dislocation - dislocation interactions (junctions, annihilation, jogs)
- ❑ Density increases with deformation (strain hardening)



Average forest dislocation spacing

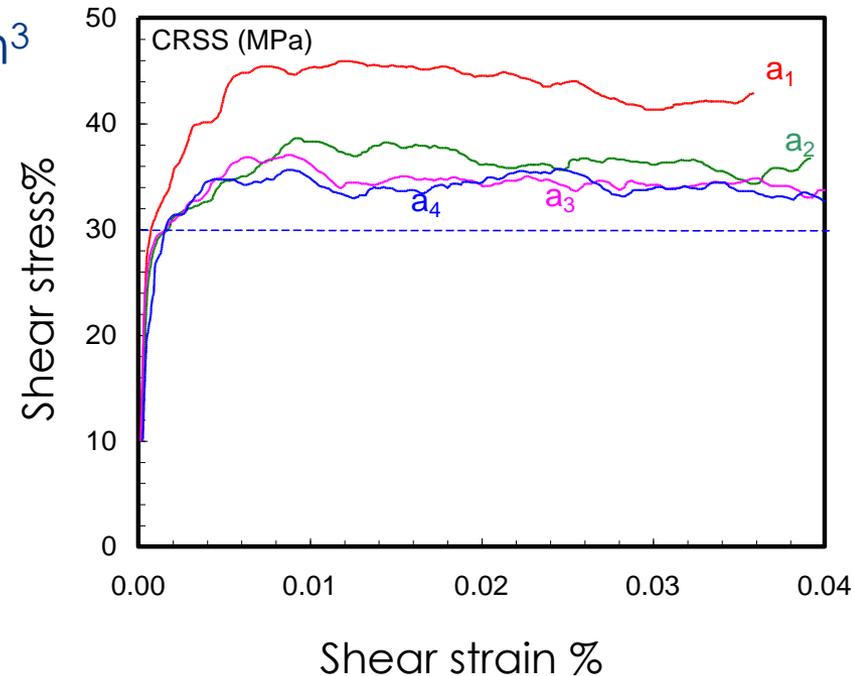
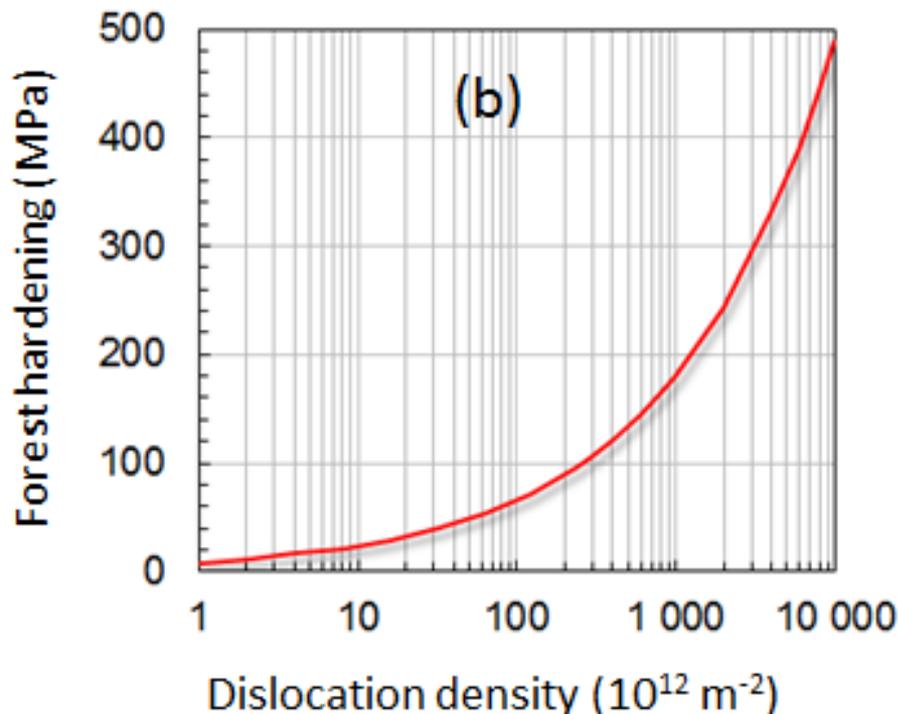
$$l = \frac{1}{\sqrt{\rho_{\text{forêt}}}}$$

Strengthening is given by :
$$\Delta\tau_{\text{forêt}} = \alpha \frac{\mu b}{l} = \alpha \mu b \sqrt{\rho_{\text{forêt}}}$$

Local obstacles: forest strengthening



- ❑ Simulation volume : $10 \times 10 \times 10 \mu\text{m}^3$
- ❑ Forest density : 10^{12}m^{-2}
- ❑ DD simulations at constant strain rate

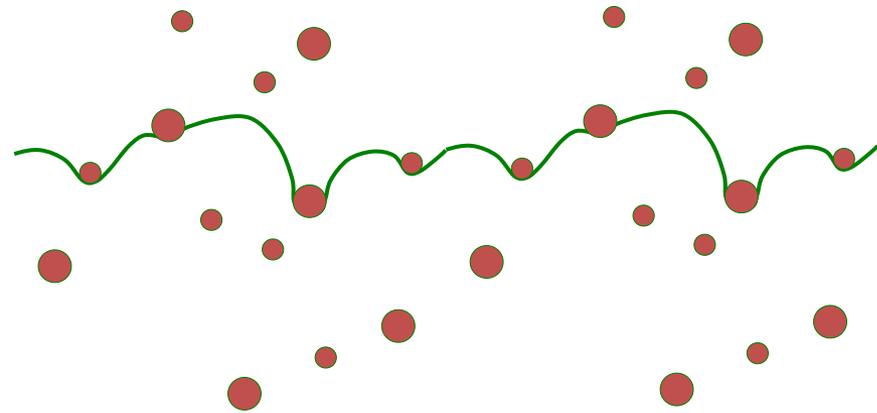


$$\Delta \tau_{forest} = \mu b \sqrt{a_{forest} \rho_{forest}}$$

Local obstacles: forest hardening in FCC



	A2	A3	A6	B2	B4	B5	C1	C3	C5	D1	D4	D6
A2	a^*	a^*	a^*	a_{col}	$a_{glissile}$	$a_{glissile}$	a_{Hirth}	$a_{glissile}$	a_{Lomer}	a_{Hirth}	a_{Lomer}	$a_{glissile}$
A3		a^*	a^*	$a_{glissile}$	a_{Hirth}	a_{Lomer}	$a_{glissile}$	a_{col}	$a_{glissile}$	a_{Lomer}	a_{Hirth}	$a_{glissile}$
A6			a^*	$a_{glissile}$	a_{Lomer}	a_{Hirth}	a_{Lomer}	$a_{glissile}$	a_{Hirth}	$a_{glissile}$	$a_{glissile}$	a_{col}
B2				a^*	a^*	a^*	a_{Hirth}	a_{Lomer}	$a_{glissile}$	a_{Hirth}	$a_{glissile}$	a_{Lomer}
B4					a^*	a^*	a_{Lomer}	a_{Hirth}	$a_{glissile}$	$a_{glissile}$	a_{col}	$a_{glissile}$
B5						a^*	$a_{glissile}$	$a_{glissile}$	a_{col}	a_{Lomer}	$a_{glissile}$	a_{Hirth}
C1							a^*	a^*	a^*	a_{col}	$a_{glissile}$	$a_{glissile}$
C3								a^*	a^*	$a_{glissile}$	a_{Hirth}	a_{Lomer}
C5									a^*	$a_{glissile}$	a_{Lomer}	a_{Hirth}
D1	Example: FCC crystallographic structure									a^*	a^*	a^*
D4											a^*	a^*
D6												a^*



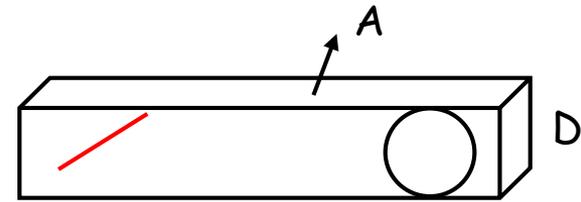
For Orowan precipitates
(impenetrable precipitates)

$$\Delta\tau_{Orowan}^s = \left(\frac{\ln 2D/b}{\ln l/b} \right)^{\frac{3}{2}} \frac{\mu b}{2\pi l} \ln \left(\frac{l}{b} \right)$$

For shearable precipitates

$$\Delta\tau_{prc}^s = \left(\frac{\Omega_{prc}}{\Omega_{\infty}} \frac{\ln 2D/b}{\ln l/b} \right)^{\frac{3}{2}} \frac{\mu b}{2\pi l} \ln \left(\frac{l}{b} \right)$$

Precipitates of size D and density C



Number of encountered precipitates

$$n = DAC$$

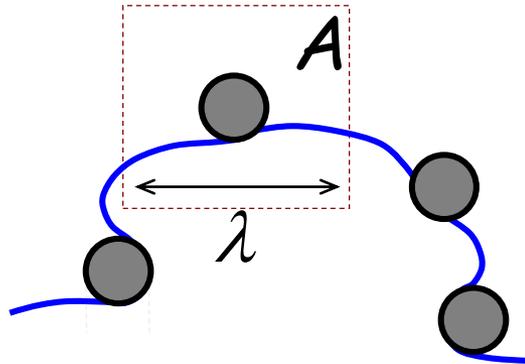
Free planner spacing

$$l = \frac{1}{\sqrt{DC}}$$

$$\Delta\tau_{prc}^s = \mu b \sqrt{a_{prc} D_{prc} C_{prc}}$$

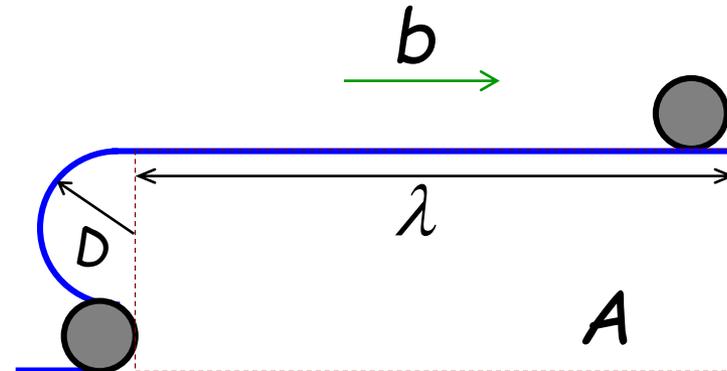
[Monnet, Acta materialia, 2015]

Random distribution of obstacles with planner density ρ



$$l = \frac{1}{\sqrt{DC}}$$

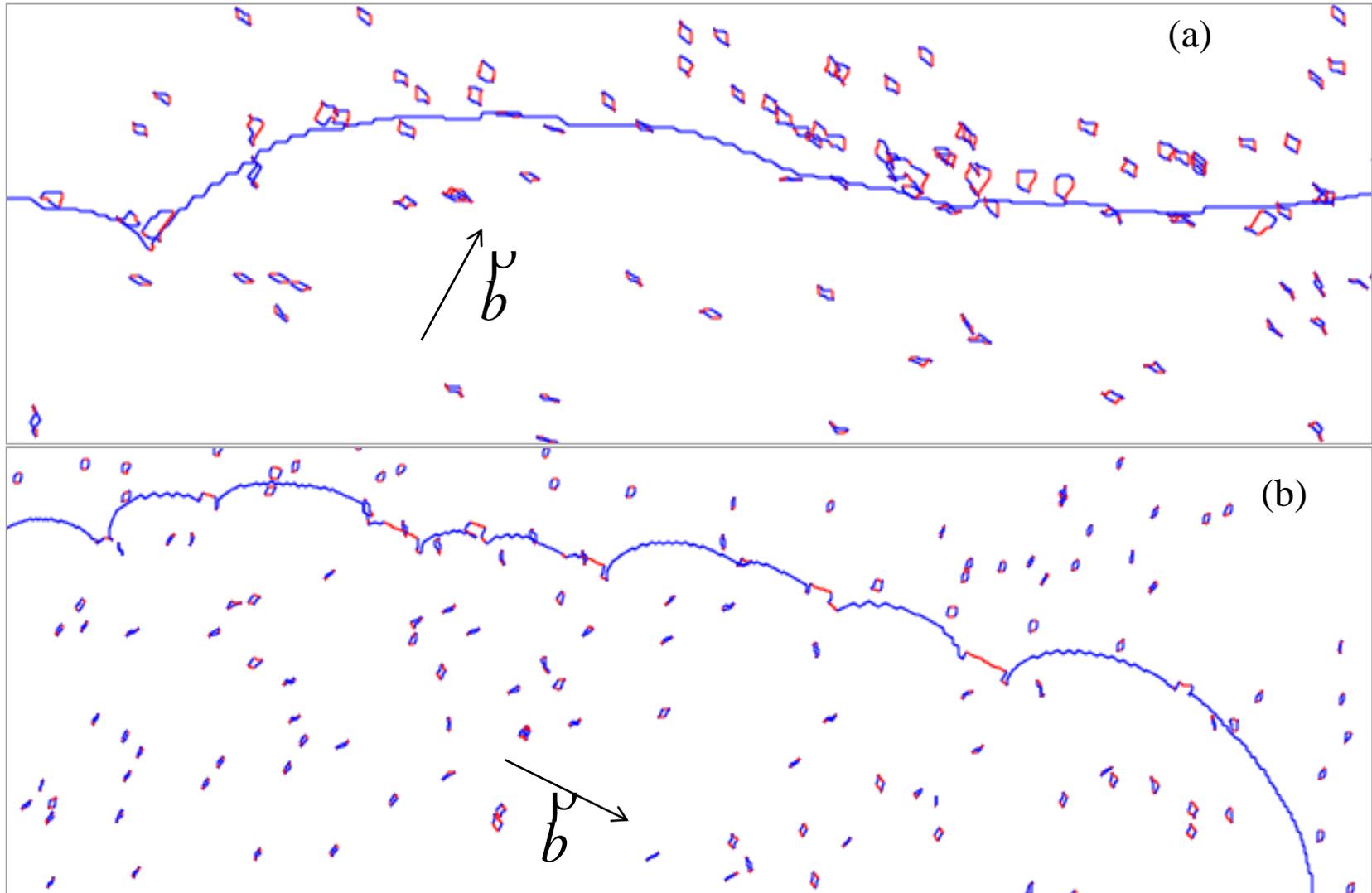
$$\tau_{obs}^s = \alpha^s \mu b \sqrt{DC}$$



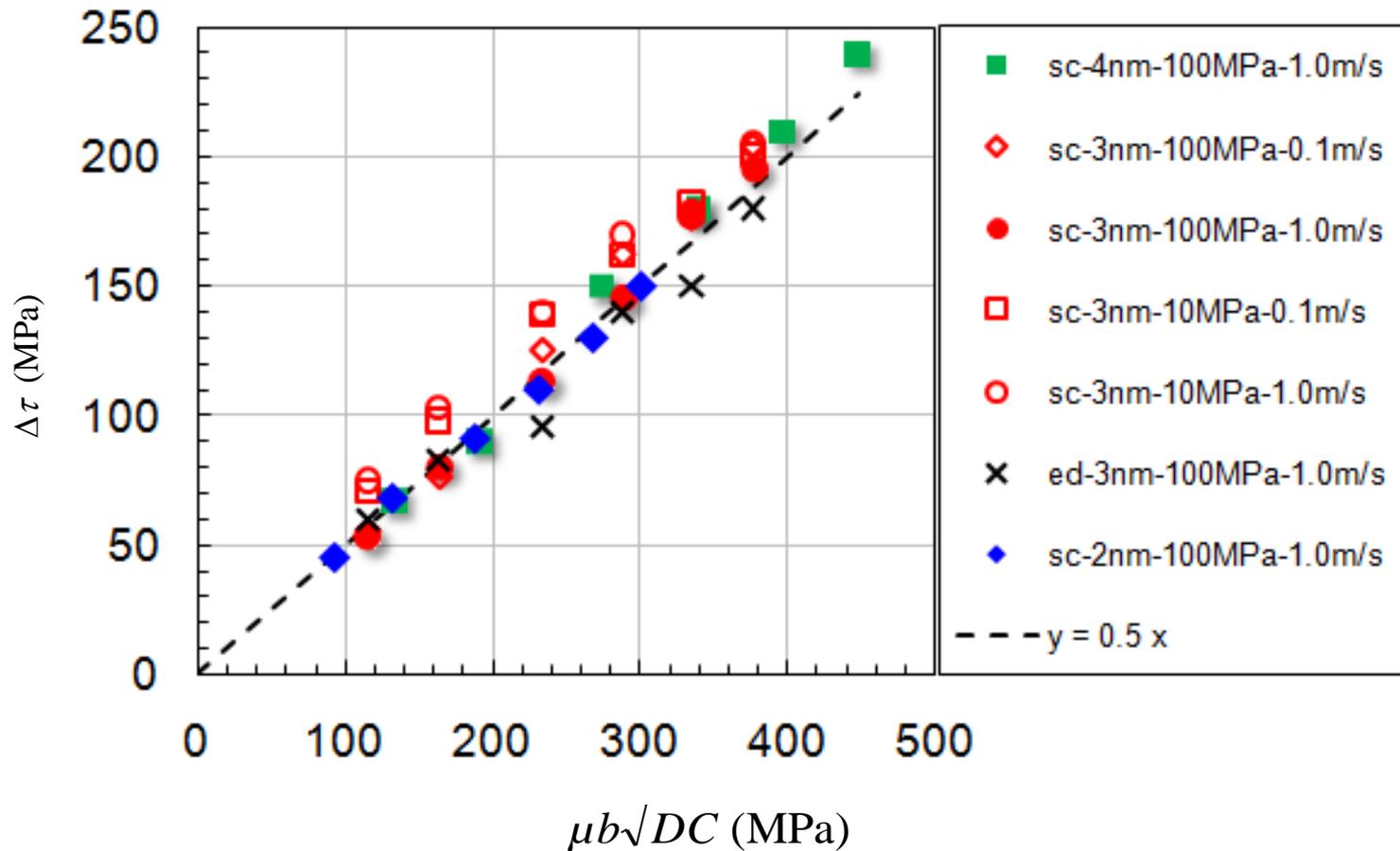
$$l = \frac{1}{(D + 2R)\rho}$$

$$\tau_{obs}^s = \begin{cases} 0 \\ \alpha^s \mu b \sqrt{DC} - \tau_{eff}^s \end{cases}$$

Local obstacles: dislocation loops

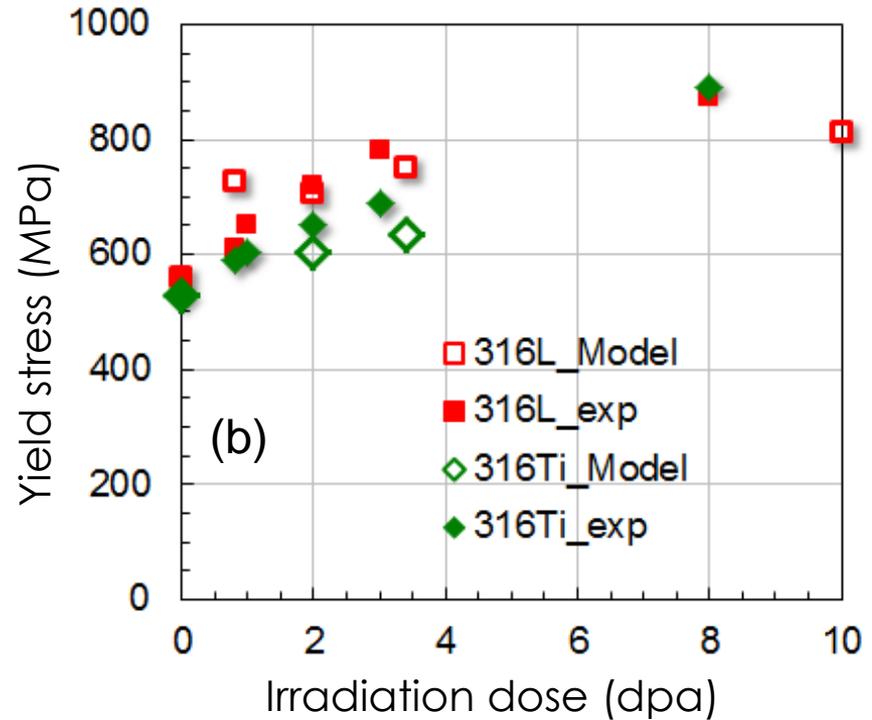
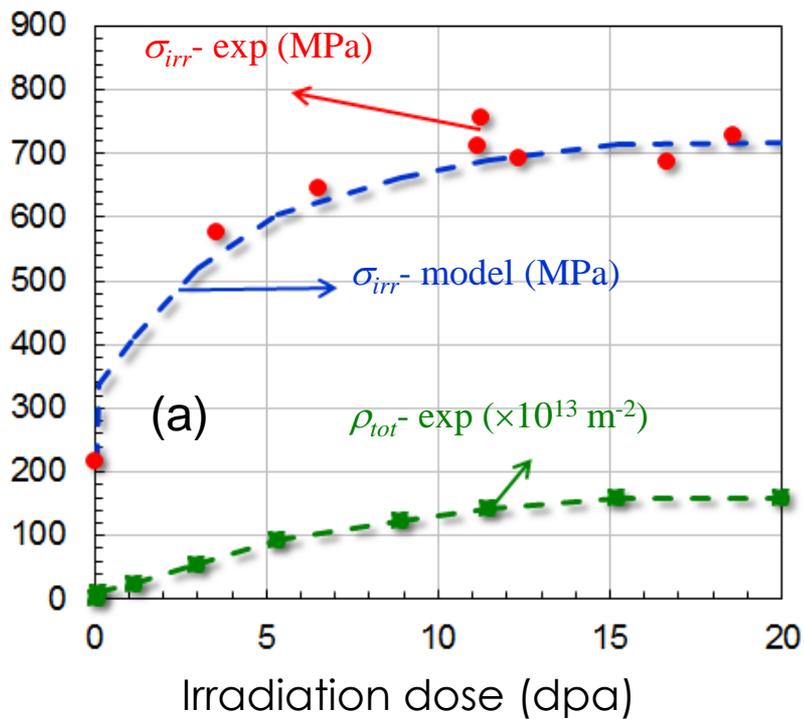


Local obstacles: dislocation loops



$$\tau_{DL} = \alpha_{forest} \mu b \sqrt{\rho_{DL}} \quad [\text{Monnet, Scripta materialia, 2015}]$$

Local obstacles: dislocation loops



$$Rp = M\tau_c = M(\tau_{ss} + \alpha\mu b\sqrt{\rho_{tot}} + \tau_{HP}) = \tau_{ss} + \alpha\mu b\sqrt{\rho_{forest} + \rho_{DL}} + \tau_{HP}$$

G. Monnet, Multiscale modeling of irradiation hardening:
Application to important nuclear materials, JNM 2018

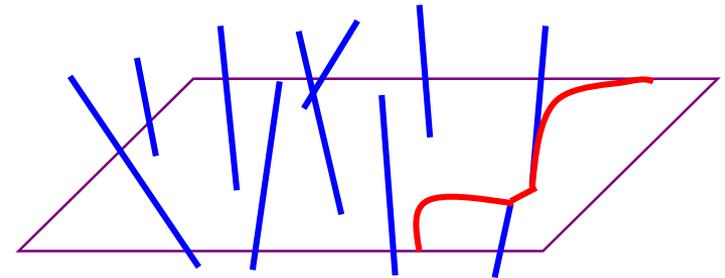
- ❑ What is a crystalline law ?
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$$d\rho^s = f(\gamma^i, T, \dots, \text{etc.})$$

- Density on system s varies with slip on system s

$$d\rho^s = f(\gamma^s)$$

- A fraction of mobile dislocations stops at the forest which leads to a generation of new dislocation



- Assuming a dislocation of length L to stop after displacement over l , the accommodated shear strain:

- The corresponding stored density:

$$\left. \begin{aligned} d\gamma^s &= \frac{\lambda L b}{V} \\ d\rho^s &= \frac{L}{V} \end{aligned} \right\} d\rho^s = \frac{d\gamma^s}{b\lambda}$$

- λ is called free mean path of dislocations

- λ scales with average spacing of the forest density

$$\lambda \propto \frac{1}{\sqrt{\rho_{forest}^s}}$$

- λ must decrease with the forest strength

$$\lambda \propto \frac{1}{\alpha \sqrt{\rho_{forest}^s}}$$

- λ can be given as:

$$\lambda \propto \frac{K}{\alpha \sqrt{\rho_{forest}^s}}$$

- λ obeys the harmonic sum for superposition

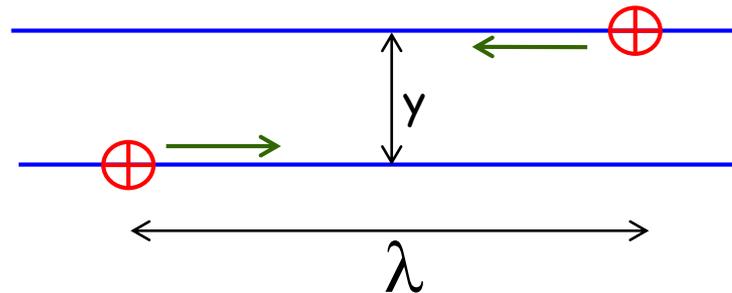
$$\frac{1}{\lambda} = \frac{1}{\lambda_1} + \frac{1}{\lambda_2} \propto \frac{\alpha_1 \sqrt{\rho^1}}{K_1} + \frac{\alpha_2 \sqrt{\rho^2}}{K_2}$$

- the total storage rate

$$d\rho^s = \frac{d\gamma^s}{b\lambda} = \frac{d\gamma^s}{b} \left(\sum_{i=1}^N \frac{\alpha_i \sqrt{\rho^i}}{K_i} \right)$$

$$\rho^s = \frac{1}{b} \left(\sum_{i=1}^N \frac{\sqrt{a^{is} \rho^i}}{K_i} \right) \rho^s$$

- Dislocation multiplication \Rightarrow closer dislocations \Rightarrow strong attraction between dipoles
- Annihilation by: cross-slip (screw) and climb (edge) dislocations \Rightarrow recovery rate



- All dislocations gliding on slip planes closer than “y” annihilate when they get close to each other
- Dislocations of length L within a distance λ can annihilate

- Annihilated density

$$d\rho^s = -\frac{Ly\lambda\rho^s}{V}$$

- Accommodated strain

$$d\gamma^s = \frac{bL\lambda}{V}$$

$$\left. \begin{array}{l} d\rho^s = -\frac{Ly\lambda\rho^s}{V} \\ d\gamma^s = \frac{bL\lambda}{V} \end{array} \right\} d\rho^s = -\frac{y\rho^s d\gamma^s}{b}$$

$$d\rho_+^s = \frac{d\gamma^s}{b} \left(\sum_{i=1}^N \frac{\alpha_i \sqrt{\rho_i^s}}{K_i} \right) \quad d\rho_-^s = -\frac{1}{b} (y\rho^s) d\gamma^s$$



$$\mathcal{R}_{total}^s = \frac{\mathcal{R}^s}{b} \left(\sum_{i=1}^N \frac{\alpha_i \sqrt{\rho^i}}{K_i} - y\rho^s \right)$$

□ Rate equation $\dot{\epsilon} = \dot{\epsilon}_0 \left(\frac{\tau_{app}^s}{\tau_c^s} \right)^n$

□ Stress decomposition $\tau_c^s = \tau_f^s + \sqrt{\mu^2 b^2 \sum_{i=1}^N a^{si} \rho^i + \Delta \tau_i^2} + \tau_{HP}$

□ Microstructure evolution $\frac{\dot{\epsilon}}{b \dot{\epsilon}} = \frac{1}{d_{grain}} + \left(\frac{\sqrt{a^{self} \rho^s}}{K_{self}} + \frac{\alpha^s \sqrt{\rho_{obs}^s}}{K_{obs}} - y \rho^s \right)$

$$\dot{C}_{DL}^s = -\lambda_{DL} \frac{D_{DL}^s}{b} C_{DL}^s |\dot{\epsilon}|$$

$$\dot{C}_{DL}^s = -\lambda_{DL} \frac{D_{DL}^s}{b} C_{DL}^s |\dot{\epsilon}|$$

Summary of the crystalline law

□ rate equation $\frac{1}{\dot{\gamma}^s} = \frac{1}{\dot{\gamma}_{drag}^s} + \frac{1}{\dot{\gamma}_{friction}^s}$

$$\dot{\gamma}_{drag}^s = \dot{\gamma}_o^s \left(\frac{\tau_{app}^s}{\tau_c^s} \right)^n$$

$$\dot{\gamma}_{friction}^s = \rho_m^s b H l_{sc}^s \exp \left(-\frac{\Delta G_o}{kT} \left(1 - \sqrt{\frac{\tau_{eff}^s}{\tau_o^s}} \right) \right)$$

□ stress decomposition $\tau_{app}^s = \tau_{eff}^s + \tau_f^s + \sqrt{\tau_{self}^s{}^2 + \tau_{LT}^s{}^2} + \tau_{HP}^g$

$$\rho_{obs}^s = \sum_{j \neq s} \rho_{dis}^j + \rho_{carbide} + \rho_{SC} + \rho_{DL}$$

$$\tau_{TL}^s = \frac{\alpha^s \mu b}{\lambda^s - l_c} - \tau_{eff}^s$$

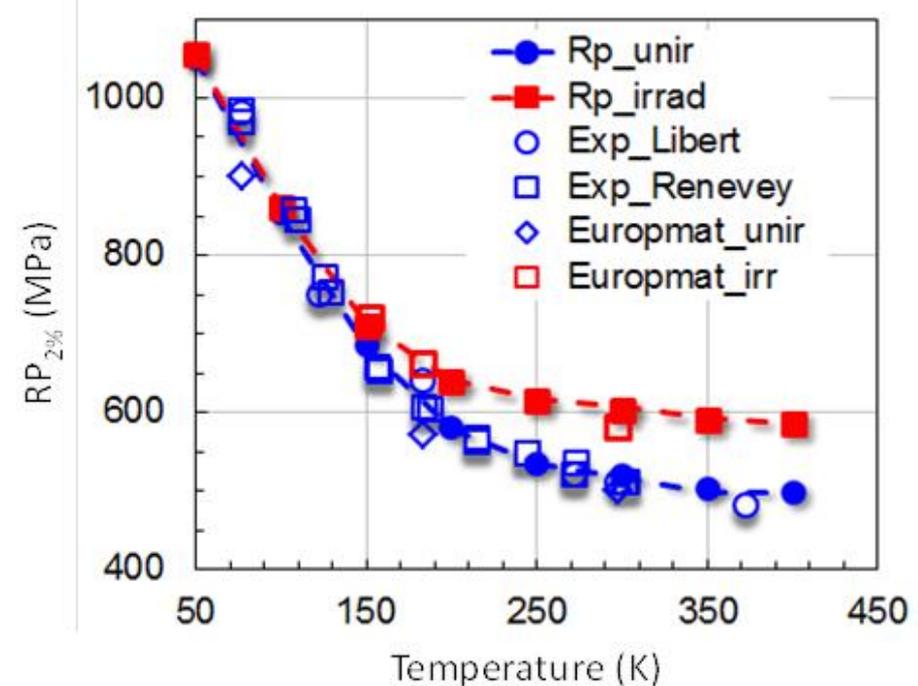
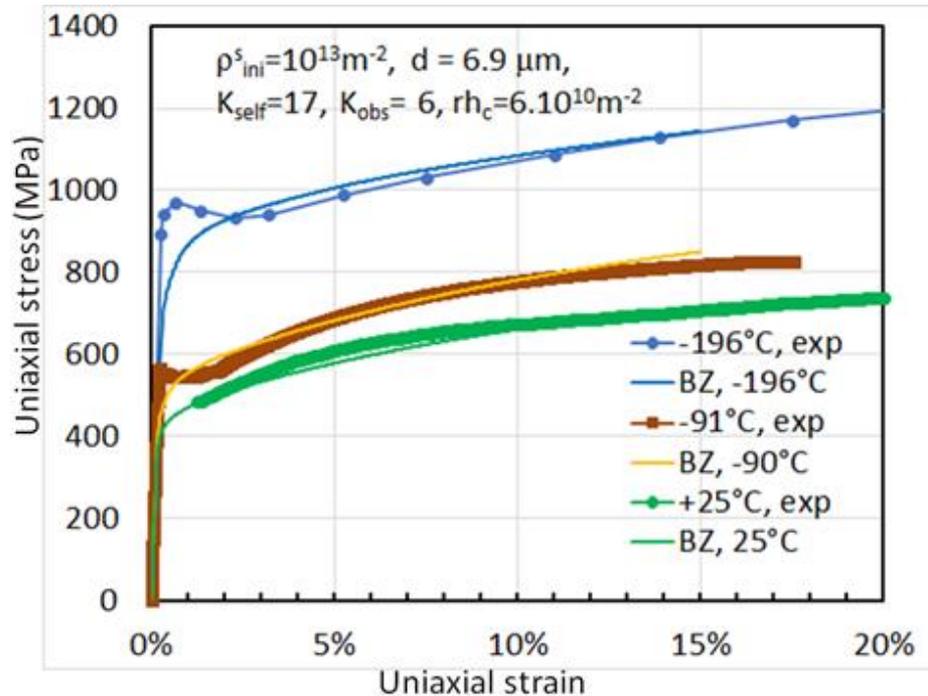
□ Microstructure evolution

$$\frac{\dot{\rho}^s}{b \dot{\gamma}^s} = \frac{1}{d_{grain}} + \left(1 - \frac{\tau_{eff}^s}{\tau_o} \right) \left(\frac{\sqrt{a^{self}} \rho^s}{K_{self}} + \frac{\alpha^s \lambda^s \rho_{obs}^s}{K_{obs}} - y \rho^s \right)$$

$$\dot{C}_{DL}^s = -\lambda_{DL} \frac{D_{DL}^s}{b} C_{DL}^s |\dot{\gamma}^s|$$

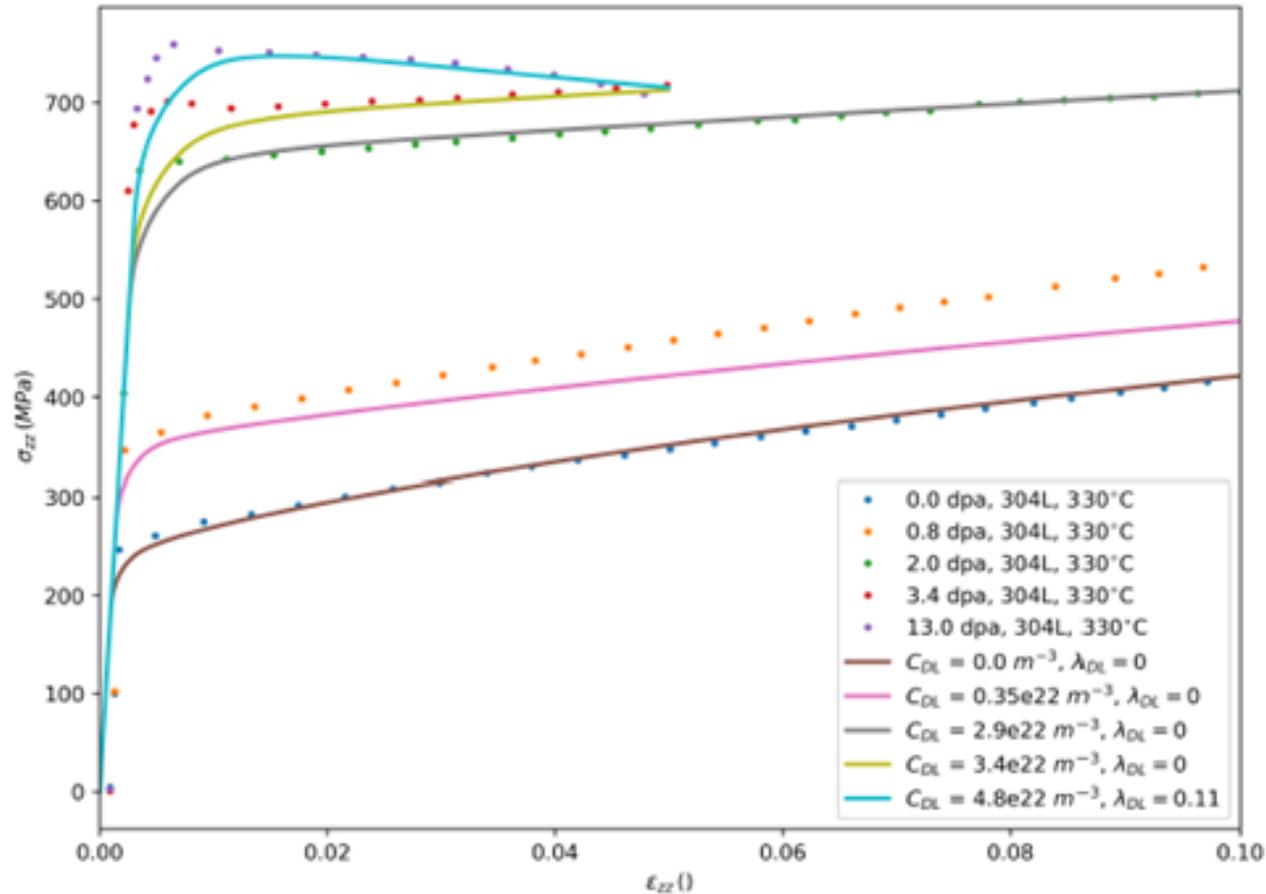
$$\dot{C}_{DL}^s = -\lambda_{DL} \frac{D_{DL}^s}{b} C_{DL}^s |\dot{\gamma}^s|$$

Validation : RPV steels



Comparison with experiment
RPV Euromaterial A

Effect of irradiation in RPV
Euromaterial A



Comparison with experiment irradiated 304L

Conclusions

- ❑ It is possible to model plastic deformation at physical basis
- ❑ Fundamental mechanisms can be treated separately
- ❑ Hardening sources cannot be added linearly

